Is it always easier to capture time-varying dynamics which resist the influence of external perturbations?

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Introduction

My work in Lancaster group together with:



Prof. Aneta Stefanovska



Dr Phil Clemson



Dr Tomislav Stankovski

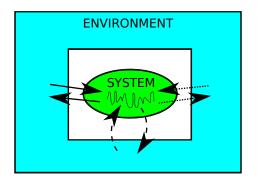
Dr Gemma Lancaster

Time-varying dynamics?

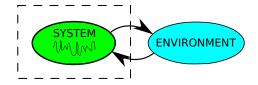
- 2 Ability to resist external perturbations?
- 3 Nonautonomous dynamical systems
- 4 Chronotaxic systems
- 5 Cardio-respiratory system

Nonautonomous (time-varying) dynamics in living (open) systems:

$$\dot{x}=f(x,t)$$



Time-varying dynamics?



Examples, just to name a few:

• in cardio-respiratory system



http://goo.gl/QXjI0v

• in cells (metabolic oscillations)



http://goo.gl/SH1tmq

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Is it always easier to capture time-varying (

Living systems:

• STRUCTURE: cells, tissues, blood vessels, organs, etc.

- often, well defined objects
- maintain their structure over long time
- certain stability

• FUNCTION: heart beat, respiration, etc.

- maintenance of stable internal environment
- certain stability

Simple mechanical analogies:

NOT STABLE function:

Oscillations with phase ϕ

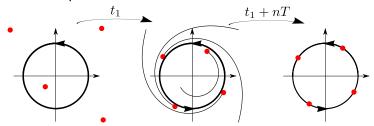
$$\dot{\phi} = \omega$$

$$\dot{\phi} = \omega + \delta \omega$$

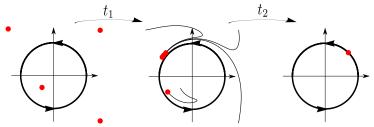
STABLE function

$$egin{array}{rcl} \dot{\phi} &=& \omega - \sin(\phi - heta) \ \dot{ heta} &=& \omega \ \psi &= \phi - heta, \ \dot{\psi} &= -\sin\psi \end{array}$$

NOT STABLE phase of oscillations



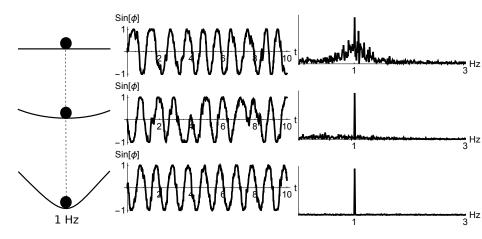
• STABLE phase of oscillations



Example: artificial heart with frequency f = 1Hz, $\dot{\phi} = g(f, \phi, t) + \xi(t)$

Ability to resist: Dynamics:

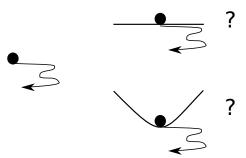
Fourier transform:



When frequency is time-dependent, when time-varying dynamics is observed...

Observed dynamics:

Does it resist perturbations?



How to quantify the ability to resist perturbations in time-varying dynamics? ... using nonautonomous systems!

Nonautonomous systems –

Skew-product flow from unidirectionally coupled differential equations,

$$\dot{\mathbf{p}} = \mathbf{f}(\mathbf{p})$$

 $\dot{\mathbf{x}} = \mathbf{g}(\mathbf{x}, \mathbf{p})$

• $\mathbf{p} \in \mathbf{R}^n, \mathbf{x} \in \mathbf{R}^m,$

- p drives the nonautonomous system x
- \mathbf{x} is dependent on t, t_0 and \mathbf{x}_0 .

Thus a system is described by:

$$\frac{d\mathbf{x}}{dt} = \tilde{\mathbf{g}}(\mathbf{x}, t)$$

Kloeden and Rasmussen, Nonautonomous Dynamical Systems (2011)

In polar coordinates (r, φ) System **p**: $\dot{r}_{\mathbf{p}} = 0$; $\dot{\varphi}_{\mathbf{p}}(t) = \omega_{\mathbf{p}}(t)$; System **x**: $\dot{r}_{\mathbf{x}} = \varepsilon_r r_{\mathbf{x}}(r_{\mathbf{p}} - r_{\mathbf{x}})$; $\dot{\varphi}_{\mathbf{x}}(t) = -\varepsilon \sin(\varphi_{\mathbf{x}} - \varphi_{\mathbf{p}})$; Black disk = $(r_{\mathbf{p}}, \varphi_{\mathbf{p}}(t))$. Gray arrows = velocities $\dot{\mathbf{x}}$; $\varepsilon > \omega_{\mathbf{p}}(t) > 0$.

Trajectories of x:

Time-dependent point attractor exists:

Time-dependent point attractor

Time-dependent point attractor (driven steady state) $\mathbf{x}^{A}(t)$.

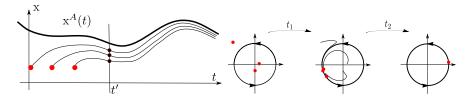
 $\mathbf{x}^{A}(t)$ satisfies mathematical conditions of pullback attraction (1), forward attraction (2) and invariance (3):

$$\lim_{t_0 \to -\infty} \mathbf{x}(t, t_0, \mathbf{x}_0) = \mathbf{x}^{\mathcal{A}}(t)$$
(1)

$$\lim_{t\to\infty} \mathbf{x}(t,t_0,\mathbf{x}_0) = \mathbf{x}^{\mathcal{A}}(t)$$
(2)

$$\mathbf{x}(t, t_0, \mathbf{x}^A(t_0)) = \mathbf{x}^A(t)$$
(3)

Kloeden and Rasmussen, Nonautonomous Dynamical Systems (2011)



Time-dependent point attractor in phase and amplitude dynamics:

We therefore define a **new class** of nonautonomous oscillators: **chronotaxic systems** (from *chronos* – time and *taxis* – order).

Suprunenko, Clemson and Stefanovska, PRL (2013); PRE (2014)

Chronotaxic Systems: Defining concepts

1) Nonautonomous systems.

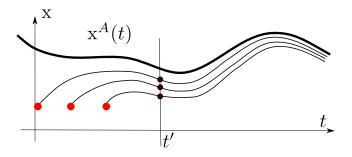
$$\frac{d\mathbf{x}}{dt} = \mathbf{g}(\mathbf{x}, t); \quad (4)$$

2) Time-dependent point attractor (driven steady state) $\mathbf{x}^{A}(t)$.

$$\lim_{t_0\to -\infty} \mathbf{x}(t, t_0, \mathbf{x}_0) = \mathbf{x}^A(t);$$
(5)

$$\lim_{t\to\infty} \mathbf{x}(t, t_0, \mathbf{x}_0) = \mathbf{x}^{\mathcal{A}}(t);$$
(6)

$$\mathbf{x}(t, t_0, \mathbf{x}^{\mathcal{A}}(t_0)) = \mathbf{x}^{\mathcal{A}}(t).$$
(7)



P. E. Kloeden and M. Rasmussen, Nonautonomous Dynamical Systems (American Mathematical Soc., 2011)

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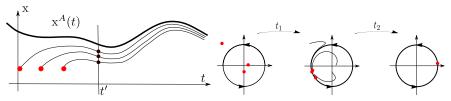
Chronotaxic systems: Definition

Chronotaxic systems are nonautonomous oscillatory systems $\mathbf{x} = \mathbf{g}(\mathbf{x}, t)$ with solutions $\mathbf{x}(t, t_0, \mathbf{x}_0)$ and with a **time-dependent point attractor (driven steady state)** $\mathbf{x}^A(t)$ which satisfies mathematical conditions of pullback attraction, forward attraction and invariance:

$$\lim_{t_0 \to -\infty} \mathbf{x}(t, t_0, \mathbf{x}_0) = \mathbf{x}^{A}(t)$$
$$\lim_{t \to \infty} \mathbf{x}(t, t_0, \mathbf{x}_0) = \mathbf{x}^{A}(t)$$
$$\mathbf{x}(t, t_0, \mathbf{x}^{A}(t_0)) = \mathbf{x}^{A}(t)$$

Additional requirement: attraction at all times, deviations $\delta \mathbf{x}^{A}(t) = \mathbf{x}(t) - \mathbf{x}^{A}(t)$ from point attractor $\mathbf{x}^{A}(t)$ in unperturbed system can only decay

$$\frac{d}{dt}|\delta \mathbf{x}^{A}|^{2} = 2\delta \mathbf{x}^{A T} J(\mathbf{x}^{A}, t)\delta \mathbf{x} < 0 \quad \text{and} \quad J(\mathbf{x}, t) = \frac{1}{2} \left(\frac{\partial \mathbf{g}}{\partial \mathbf{x}} + \frac{\partial \mathbf{g}}{\partial \mathbf{x}^{T}} \right)$$



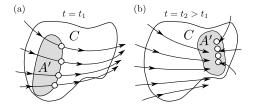
Chronotaxic systems: Alternative definition

Alternative formulation is based on the contraction theory (*W. Lohmiller and J.-J. Slotine (1998)*). The explicit knowledge of x^A(t) is not required.

A chronotaxic system is an oscillatory nonautonomous dynamical system $\dot{\mathbf{x}} = \mathbf{g}(\mathbf{x}, t)$ which has a contraction region *C* in the phase space, *C* is determined by

$$\forall \mathbf{x} \in C \; \exists \beta > 0, \forall t : \qquad \frac{1}{2} \left(\frac{\partial \mathbf{g}(\mathbf{x}, t)}{\partial \mathbf{x}} + \frac{\partial \mathbf{g}(\mathbf{x}, t)}{\partial \mathbf{x}}^T \right) \leq -\beta I < 0,$$

C contains a finite non-zero area $A' \subset C$ such that states of a system inside A' do not leave A', i.e. $\forall t_0 < t$, $\forall \mathbf{x}_0 \in A'(t_0)$, $\mathbf{x}(t, t_0, \mathbf{x}_0) \in A'(t)$.



Y. S., Stefanovska A., PRE (2014)

Chronotaxic systems: Phase oscillators

Phase α of the phase oscillator which is

conventional

chronotaxic

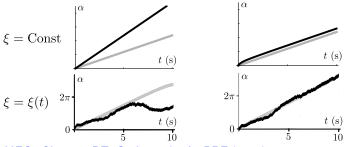
does not resist external perturbations ξ , as shown in this column:

is able to resist external perturbations ξ , as shown in this column:

$$\dot{\alpha} = \omega_0 + \xi$$

$$\left. \frac{\partial}{\partial lpha} g_{lpha}(e^{i lpha_{x}},t) \right|_{lpha=lpha^{A}(t)} < 0.$$

 $\dot{\alpha} = q_{x} (e^{i\alpha_{x}} t) + \xi$



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Chronotaxic systems: Poincare oscillator

Nonautonomous (driven) Poincare oscillator:

$$\dot{x} = \varepsilon_{\Gamma} \left(r_{\rho} - \sqrt{x^2 + y^2} \right) x - \omega_0 y - \varepsilon_A(t) \left(x - r_{\rho} \cos \varphi_{\rho}(t) \right) + \xi_1(t),$$

$$\dot{y} = \varepsilon_{\Gamma} \left(r_{\rho} - \sqrt{x^2 + y^2} \right) y + \omega_0 x - \varepsilon_A(t) \left(y - r_{\rho} \sin \varphi_{\rho}(t) \right) + \xi_2(t).$$

- nonautonomous system:
 p is defined by polar coordinates r_p, φ_p(t);
 x is defined by x, y;
- $\varepsilon_A(t)$ determines coupling;
- perturbations are defined by $\xi_1(t), \xi_2(t)$;
- phase and amplitude dynamics are coupled.

Y.F.S., Stefanovska A., PRE (2014)

Nonautonomous (driven) Poincare oscillator: $\dot{\varphi}_{p}(t) = \omega_{p}(t) = \text{const};$

$$\dot{x} = \varepsilon_{\Gamma} \left(r_{\rho} - \sqrt{x^2 + y^2} \right) x - \omega_0 y - \varepsilon_A(t) \left(x - r_{\rho} \cos \varphi_{\rho}(t) \right) + \xi_1(t),$$

$$\dot{y} = \varepsilon_{\Gamma} \left(r_{\rho} - \sqrt{x^2 + y^2} \right) y + \omega_0 x - \varepsilon_A(t) \left(y - r_{\rho} \sin \varphi_{\rho}(t) \right) + \xi_2(t).$$

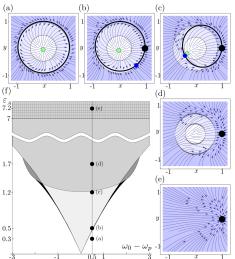


Figure shows:

 snapshots of phase portrait of this system when it is:

(a) non-chronotaxic;

(b),(c),(d),(e) - chronotaxic (time-dependent point attractor is shown as a black disk, the green point is an unstable node, the blue point is a saddle point);

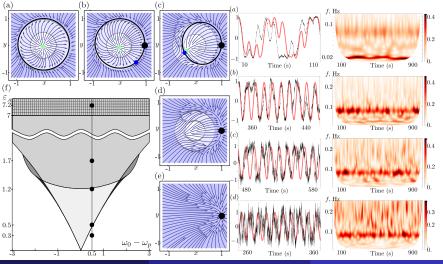
 (f) regions of different chronotaxic dynamics.

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- Parameters: white Gaussian noise $\langle \xi_j(t) \rangle = 0$, $\langle \xi_i(t) \xi_j(t') \rangle = \sigma^2 \delta_{ij} \delta(t t')$, $\omega_p/(2\pi) \approx 0.08$ Hz, $\varepsilon_r = 7$, $r_p = 1$; (a) $\varepsilon_A = 0.3$, $\sigma = 0.1$; (b) $\varepsilon_A = 0.47$, $\sigma = 0.3$ (c) $\varepsilon_A = 0.47$, $\sigma = 0.6$ (d) $\varepsilon_A = 0.9$, $\sigma = 1.2$.
- In chronotaxic systems (b),(c),(d) the frequency ω_p/(2π) determines a frequency of 0.08Hz which resists external perturbations, as shown in wavelet transforms;

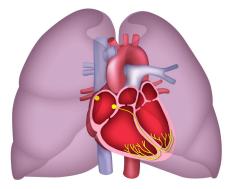


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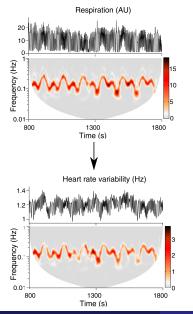
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Cardio-respiratory system



Cardio-respiratory system



Y.F.S., Clemson P.T., Stefanovska A., PRL **111**, 024101 (2013) Clemson P.T., Y.F.S., Stefanovska A., PRE (2014).

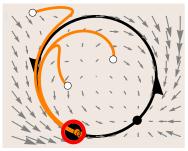
- (1) The paced respiration drives HRV;
- We are looking for a new feature a stability of a time-varying dynamics;
- (3) Stability in time-varying HRV is found using a single time-series, and checked using both time-series;
- (4) Alterations in chronotaxicity may help to distinguish normal and altered states.



Summary

Chronotaxicity as an ability to resist external perturbations:

- helps to capture time-variability;
- is a quantifiable characteristics;
- may help to identify subsystems which maintain their function, and to clarify functional organization within a living system.



Thank you for your attention!